



A feasibility study of the time reversal violation test based on polarization of annihilation photons from the decay of ortho-Positronium with the J-PET detector

J. Raj¹ · A. Gajos¹ · C. Curceanu² · E. Czerwiński¹ · K. Dulski¹ · M. Gorgol³ · N. Gupta-Sharma¹ · B. C. Hiesmayr⁴ · B. Jasińska³ · K. Kacprzak¹ · Ł. Kapłan¹ · D. Kisielewska¹ · K. Klimaszewski⁵ · G. Korcyl¹ · P. Kowalski⁵ · T. Kozik¹ · N. Krawczyk¹ · W. Krzemień⁵ · E. Kubicz¹ · M. Mohammed⁶ · Sz. Niedźwiecki¹ · M. Pałka¹ · M. Pawlik-Niedźwiecka¹ · L. Raczyński⁵ · K. Rakoczy¹ · Z. Rudy¹ · S. Sharma¹ · Shivani¹ · R. Y. Shopa⁵ · M. Silarski¹ · M. Skurzok¹ · W. Wiślicki⁵ · B. Zgardzińska³ · P. Moskal¹

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Abstract

The Jagiellonian Positron Emission Tomograph (J-PET) is a novel device being developed at Jagiellonian University in Krakow, Poland based on organic scintillators. J-PET is an axially symmetric (M. Mohammed) has been affiliated to affiliation 6. Please check if correct, otherwise, provide the correct affiliation. and high acceptance scanner that can be used as a multi-purpose detector system. It is well suited to pursue tests of discrete symmetries in decays of positronium in addition to medical imaging. J-PET enables the measurement of both momenta and the polarization vectors of annihilation photons. The latter is a unique feature of the J-PET detector which allows the study of time reversal symmetry violation operator which can be constructed solely from the annihilation photons momenta before and after the scattering in the detector.

Keywords J-PET detector · Time reversal symmetry · Ortho-positronium · Discrete symmetries · Momenta · Polarization · Annihilation photons

1 Introduction

Testing whether Nature is invariant under time reversal is one of the most curious challenges in physics. Although, one can reverse the direction of motion in space, one cannot

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✉ J. Raj
juhi.raj@doctoral.uj.edu.pl

Extended author information available on the last page of the article.

Table 1 Discrete symmetry odd-operators using spin orientation of the o-Ps and momentum directions of the annihilation photons

Operator	C	P	T	CP	CPT
$S \cdot k_1$	+	−	+	−	−
$S \cdot (k_1 \times k_2)$	+	+	−	+	−
$(S \cdot k_1) \cdot (S \cdot (k_1 \times k_2))$	+	−	−	−	+

Table 2 Discrete symmetry odd-operators using spin orientation of the o-Ps as well as polarization and momentum directions of the annihilation photons

Operator	C	P	T	CP	CPT
$\epsilon_1 \cdot k_2$	+	−	−	−	+
$S \cdot \epsilon_1$	+	+	−	+	−
$S \cdot (k_2 \times \epsilon_2)$	+	−	+	−	−

reverse the direction of elapsing time. Therefore, in order to test the time reversal violation, we investigate the expectation values for the T-odd-symmetry-operators. To date, time reversal symmetry violation has not been observed in purely leptonic systems. The best experimental upper limits for CP and CPT (C-Charge Conjugation, P-Parity and T-Time) symmetries violation in positronium decay are set to 0.3×10^{-3} [1–3]. According to the Standard Model predictions, photon-photon interaction or weak interaction can mimic the symmetry violation in the order of 10^{-9} (photon-photon interaction) and 10^{-13} (weak interactions), respectively [4–7]. There is about 7 orders of magnitude difference between the present experimental upper limit and the standard model predictions. So far, discrete symmetries were proposed to be tested with the ortho-positronium (o-Ps) system by determining the non-zero expectation values of the operators constructed from the spin (S) of the o-Ps and the momentum vectors of the annihilation photons k_1 , k_2 and k_3 as mentioned in Table 1.

In this article we explore feasibility of applying a new method of studying the fundamental discrete symmetry violation in a purely leptonic system proposed in Ref. [8]. The novelty of the method is constituted by using of the discrete symmetry odd-operators constructed with the polarization (ϵ_i) of the annihilation photons originating from the decay of o-Ps atoms, as listed in Table 2 [8]. The observation of non-zero expectation values of these operators would imply non-invariance of these symmetries for which the given operator is odd (marked “−” in Tables 1 and 2).

2 Jagiellonian - positron emission tomograph

Positron emission tomography (PET) is a non-invasive technique used in the diagnosis of various types of tumors at the cellular level. All commercially available PET-scanners utilize relatively expensive crystal detectors for the detection of annihilation photons [16, 17].

The Jagiellonian - Positron Emission Tomograph (J-PET) is the first PET-scanner constructed using plastic scintillator strips to make the tomograph cost effective and portable [9, 10, 13]. One of the unique features of the J-PET detector is its ability to measure polarization of the annihilation photons. The J-PET detector consists of 192 plastic scintillator strips (EJ-230) of dimensions $500 \times 19 \times 7 \text{ mm}^3$ each, forming three concentric layers (48 modules on radius 425 mm, 48 modules on radius 467.5 mm and 96 modules on radius 575 mm) (Fig. 1)[13]. Each scintillator in the J-PET scanner is optically connected at two ends to Hamamatsu R9800 vacuum tube photomultipliers [10]. This results in 384 analog channels

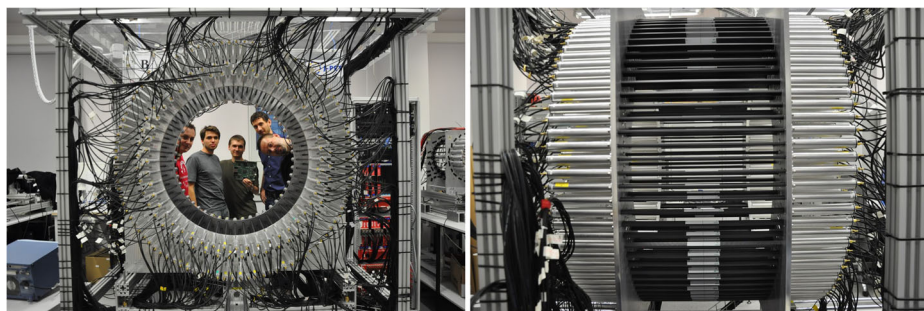


Fig. 1 (Left) Front view of the J-PET detector with a few members of the J-PET group. (Right) Side view of the J-PET detector, where, the black strips are the scintillators wrapped in light protective foil. The silver tubes are housings of photomultipliers which are connected at either ends of the scintillator

that are processed by a fully equipped trigger-less Data Acquisition System (DAQ) together with the readout mechanism [14, 18]. The readout chain consists of Front-End Electronics (FEE) such as Time-to-Digital or Analog-to-Digital Converters (TDCs or ADCs), data collectors and storage [14, 18].

The J-PET collaboration has developed a dedicated J-PET analysis framework which is a highly flexible, ROOT-based software package which aids the reconstruction and calibration procedures for the tomograph. The J-PET analysis framework stands as a backbone to analyze the collected data to produce the results presented in this article [15].

The J-PET detector, together with the trigger-less DAQ system constitutes an efficient photon detector with high timing properties [13]. This allows us to investigate the fundamental discrete symmetries in the purely leptonic sector [8].

3 Measurement technique

3.1 Production of positronium

In order to produce positronium atoms, a point-like ^{22}Na source is placed in the center of the detector and is surrounded with XAD-4 porous polymer [19]. The porous polymer enhances the production of positronium atoms.

The positrons emitted from the source ($^{22}\text{Na} \rightarrow ^{22}\text{Ne}^* + e^+ + \nu_e$), interact with the electrons in the XAD-4 porous polymer producing a meta-stable triplet state, ortho-Positronium (o-Ps) which predominantly decays into three photons due to charge conjugation symmetry conservation [7]: $e^+ + e^- \rightarrow o\text{-Ps} \rightarrow 3\gamma$. Meanwhile, the excited $^{22}\text{Ne}^*$ de-excites emitting gamma quantum with an energy of 1274 keV ($^{22}\text{Ne}^* \rightarrow ^{22}\text{Ne} + \gamma_{(1274\text{ keV})}$).

3.2 Selection of o-Ps decay events

For the study of time reversal symmetry in o-Ps we are interested in events with five interactions which are due to: one de-excitation photon, three primary annihilation photons, and one secondary scattered photon (k'_1), as indicated in Fig. 2. The signal (hit) multiplicity distribution presented in the left panel of Fig. 3 shows that the requirement of five hits in one event reduces the measured data sample by a factor of about 10^3 . The de-excitation photon is identified using the time-over-threshold (TOT) measurement [20] which is related to the

Fig. 2 Schematic of the J-PET detector with three distinguishable layers of scintillators at blue, red and yellow concentric cylinders. A positron source (red) placed in the center, covered in XAD-4 porous polymer (blue). The superimposed arrows indicate gamma photon originating from the de-excitation of $^{22}\text{Ne}^*$ (γ), annihilation photons from ortho-positronium decay (k_1 , k_2 and k_3), and scattered photon (k'_1)

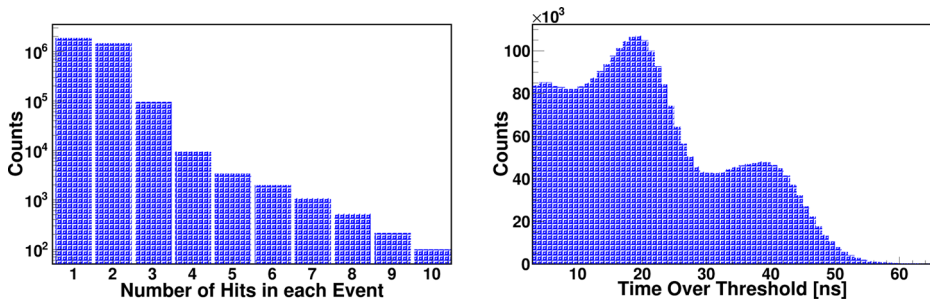
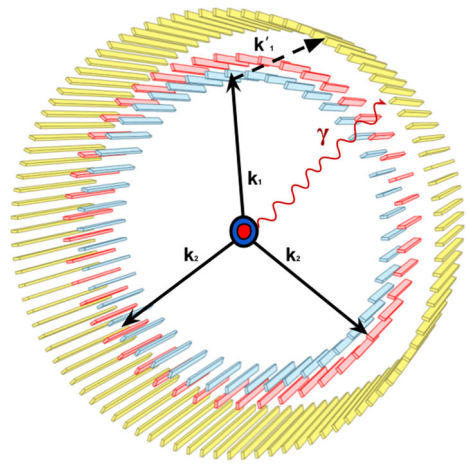


Fig. 3 (Left) The distribution of number of gamma interactions recorded in each event. (Right) Time over threshold of photomultiplier signals, used as a measure of the energy deposited in the scintillator strips

energy deposited in the scintillator. Right panel of Fig. 3 presents TOT distribution where one can clearly recognize Compton spectra from 511 keV and 1274 keV gamma photons. The de-excitation photon (1274 keV) may be selected with the efficiency of about 66% when requiring TOT larger than 30 ns.

After selection of the de-excitation photon, we need to identify events in which annihilation photons originate from the decay of ortho-positronium, and relate the primary annihilation photon to its corresponding secondary scattered photon.

From kinematics we infer that the sum of the two smallest relative azimuthal angles between the registered annihilation photons for $o\text{-Ps} \rightarrow 3\gamma$ must be greater than 180° [11, 12]. Therefore, the events corresponding to the decay of $o\text{-Ps} \rightarrow 3\gamma$ lie at the right side of the band at 180° (Fig. 4(Left)) [11, 12]. The signals corresponding to the annihilation photons and secondary scattered photon may be identified after ordering the hits in time. The signal from the scattered photon should appear as the last one. As a last step we need to assign the scattered photon to its primary photon. For this purpose we introduce a parameter $\Delta_{ij} = (\delta_M - \delta_C)$, where, δ_M and δ_C are the measured time difference and calculated time of flight between the i^{th} and j^{th} interaction, respectively. Therefore, Δ_{ij} should be equal to zero in case if the j^{th} signal is due to the i^{th} scattered photon. Right panel of Fig. 4 shows an

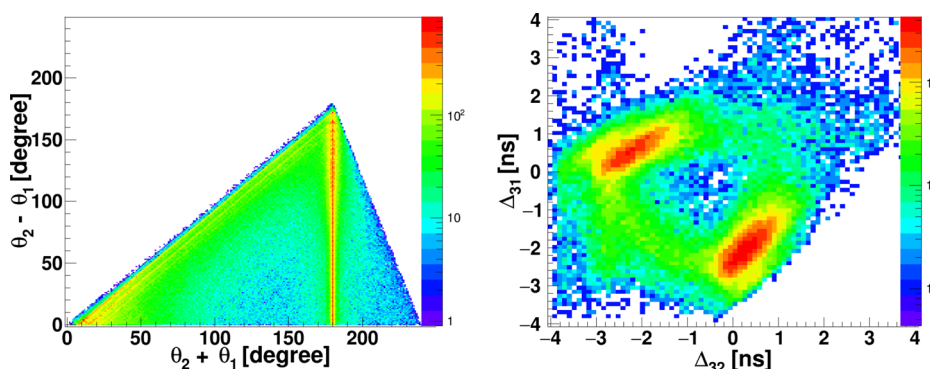


Fig. 4 (Left) Relation between the sum and difference of two smallest relative azimuthal angles (θ_1 and θ_2 (Fig. 2-Left)) of the three interacting annihilation photons. (Right) Distribution of Δ_{31} vs Δ_{32} as explained in the text

exemplary spectrum allowing to relate the primary annihilation photon to its corresponding secondary photon.

Finally, the direction of polarization (ϵ_1) of the 1st (most energetic) annihilation photon is estimated as the cross-product of the momentum of the primary and secondary photons ($k_1 \times k'_1$) [8], and is used in the determination of the expectation value of the time symmetry odd-operator $\epsilon_1 \cdot k_2$.

4 Conclusion

Positronium is an excellent laboratory enabling studies of many interesting phenomena [21] such as e.g. gravitation of antimatter [21], search for mirror photons [22, 23], quantum entanglement [24, 25] and tests of discrete symmetries in the leptonic sector [8]. So far, T-violation effects were not discovered in a purely leptonic system e.g in the positronium system. The Jagiellonian-PET scanner offers high acceptance in the measurement of the direction of polarization (ϵ) of photons simultaneously with their momentum direction (k). In reference [11] it was shown that the J-PET detector is capable of registration of $oPs \rightarrow 3\gamma$ decays with the rate ranging between 25 to 18300 events per second, depending on the configuration and amount of the used scintillator strips. In 2018 J-PET will be updated with 320 new scintillator strips leading to the registration rate of about 450 per second for the $oPs \rightarrow 3\gamma$ events. A need of registration of the additional scattered photon for the studies of $\epsilon_1 \cdot k_2$ operator will reduce this rate to about 30 events per second. Thus during one year of the data collection (achievable in practice in about 3 years) it is expected to collect about 10^9 events and hence it is expected to increase the sensitivity by about order of magnitude with respect to previous analyses based on about 10^7 events [2, 3].

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References

1. Kostelecky, V.A., Russell, N.: *Rev. Mod. Phys.* **83**, 11 (2011)
2. Yamazaki, T., et al.: *Phys. Rev. Lett.* **104**, 083401 (2010)
3. Vetter, P.A., Freedman, S.J.: *Phys. Rev. Lett.* **91**, 263401 (2003)
4. Sozzi, M.S.: *Discrete symmetries and CP violation*, Oxford University Press (2008)
5. Bernreuther, W. et al.: *Z. Phys. C* **41**, 143 (1988)
6. Arbib, B.K., et al.: *Phys. Rev. A* **37**, 3189 (1988)
7. Pokraka, A., Czarnecki, A.: *Phys. Rev.* **D96**, 093002 (2017)
8. Moskal, P., et al.: *Acta Phys. Polon. B* **47**, 509 (2016)
9. Moskal, P., et al.: *Phys. Med. Biol.* **61**, 2025–2047 (2016)
10. Moskal, P., et al.: *Nucl. Instr. Meth. A* **764**, 317 (2014)
11. Kamińska, D., et al.: *Eur. Phys. J. C* **76**, 445 (2016)
12. Gajos, A., et al.: *Nucl. Instrum. Methods A* **819**, 54 (2016)
13. Niedźwiecki, S., et al.: *Acta Phys. Pol. B* **48**, 1567 (2017)
14. Korcyl, G., et al.: *Acta Phys. Polon. B* **47**, 491 (2016)
15. Krzemień, W., et al.: *Acta Phys. Pol. A* **127**, 1491–1494 (2015)
16. Vandenberghe, S., et al.: *EJNMMI Phys.* **3**(1), 3 (2016)
17. Słomka, P., et al.: *Semin. Nucl. Med.* **46**(1), 5–19 (2016)
18. Korcyl, G., et al.: *IEEE Trans. Med. Imaging* **37**, 11 (2018)
19. Jasińska, B., et al.: *Acta Phys. Polon. B* **47**, 453 (2016)
20. Paka, M., et al.: *JINST* **12**, P08001 (2017)
21. Cassidy, D.B.: *Eur. Phys. J. D* **72**, 53 (2018)
22. Vigo, C., et al.: arXiv: [1803.05744v2](https://arxiv.org/abs/1803.05744v2) [hep-ex] (2018)
23. Gninenko, S.N., Krasnirov, N.V.: *Mod. Phys. Lett. A* **17**(26), 1713–1724 (2002)
24. Hiesmayr, B.C., Moskal, P.: *Sci. Rep.* **7**, 15349 (2017)
25. Nowakowski, M., Fierro, D.B.: *Acta Phys. Polon. B* **48**, 1955 (2017)

Affiliations

J. Raj¹ · A. Gajos¹ · C. Curceanu² · E. Czerwiński¹ · K. Dulski¹ · M. Gorgol³ · N. Gupta-Sharma¹ · B. C. Hiesmayr⁴ · B. Jasińska³ · K. Kacprzak¹ · Ł. Kapłon¹ · D. Kisiełewska¹ · K. Klimaszewski⁵ · G. Korcyl¹ · P. Kowalski⁵ · T. Kozik¹ · N. Krawczyk¹ · W. Krzemień⁵ · E. Kubicz¹ · M. Mohammed⁶ · Sz. Niedźwiecki¹ · M. Pałka¹ · M. Pawlik-Niedźwiecka¹ · L. Raczynski⁵ · K. Rakoczy¹ · Z. Rudy¹ · S. Sharma¹ · Shivani¹ · R. Y. Shopa⁵ · M. Silarski¹ · M. Skurzok¹ · W. Wiślicki⁵ · B. Zgardzińska³ · P. Moskal¹

¹ Faculty of Physics, Astronomy and Applied Computer Science, Jagiellonian University, 30-348 Cracow, Poland

² National Institute of Nuclear Physics, Laboratori Nazionali di Frascati, 00044, Frascati, Italy

³ Department of Nuclear Methods, Institute of Physics, Maria Curie-Skłodowska University, 20-031, Lublin, Poland

⁴ Faculty of Physics, University of Vienna, 1090, Vienna, Austria

⁵ Świerk Department of Complex Systems, National Centre for Nuclear Research, 05-400 Otwock, Poland

⁶ Department of Physics, College of Education for Pure Sciences, University of Mosul, Mosul, Iraq